

High performance dual-gate bilayer graphene FET with neutral channel for logic applications

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The use of bilayer graphene (BLG) is considered as an interesting alternative to circumvent the zero-bandgap problem in single-layer graphene. In fact, the application of a transverse electric field (E) in BLG allows for the breaking of the inversion symmetry and therefore the occurrence of an electronic bandgap, whose magnitude can be tuned by the strength of the applied field [1,2]. In addition to the electrostatic gating, chemical doping of the BLG top surface induces an additional electric field that enhances the vertical displacement field [3,4].

In this work, we report a novel dual-gate BLG-FET architecture with neutral channel, in which the electrostatic control of the BLG active channel is enhanced by the use of chemical functionalization by (F₄TCNQ). In this way, the BLG-FET achieves a $I_{on}/I_{off} \sim 60$ and an electron mobility up to $\sim 1600 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$. The total charge in our devices is as low as 10^{11} cm^{-2} .

Fig. 1 shows a schematic of a F₄TCNQ-doped, dual-gate BLG-FET. A top gate high-k/metal stack (Ti/Pd/Al₂O₃) is directly deposited on a BLG film covered with F₄TCNQ. The bottom gate is represented by the 300nm SiO₂/Si stack on which BLG was previously exfoliated and identified by a combination of optical microscopy and Raman spectroscopy.

The transfer characteristics I_d vs. V_{TG} (parameterized by V_{BG}) are shown in Fig. 2. The BLG top layer is “chemically gated” by the F₄TCNQ layer, which generates a permanent E perpendicular to the BLG planes, responsible for a 30–50 meV bandgap. This E yields the asymmetry between the two BLG layers and results in a significant bandgap opening. In fact, this already results in an $I_{on}/I_{off} \sim 10$ at $V_{BG}=0$ and a total channel charge density $n \sim 4 \times 10^{12} \text{ cm}^{-2}$. For $V_{BG} < 0$, n increases to $\sim 6.5 \times 10^{12} \text{ cm}^{-2}$ due to hole injection in the bottom BLG layer. At the same time, the application of a negative V_{BG} generates an E aligned in the opposite direction to that arising by the F₄TCNQ chemical gating. Therefore, the bandgap vanishes, I_{on}/I_{off} decreases (Fig. 2). For $V_{BG} > 0$, the E enhances the E generated by the F₄TCNQ layer, therefore increasing the bandgap magnitude and boosting the I_{on}/I_{off} to values up to ~ 60 . Our BLG-FETs are normally OFF ($V_{TG} \sim 0$). To switch them on, we only need to apply a $V_{TG} \sim 3 \text{ V}$. In the non-neutral channel BLG-FETs reported in literature, I_{on}/I_{off} switching occurs at very high voltages [2], thus dramatically increasing the static power needed to drive the devices. In conditions of channel charge neutrality ($n \sim 10^{11} \text{ cm}^{-2}$). The electron mobility increases from ~ 400 to $\sim 1600 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ upon restoring the channel charge neutrality from $n \sim 4 \times 10^{12} \text{ cm}^{-2}$.

In summary, we report a novel dual-gate BLG-FET architecture that significantly improves the state-of-the-art BLG devices relying on the application of an external E for bandgap opening. Our devices achieve $I_{on}/I_{off} \sim 60$

without significantly suppressing the carrier mobility, and, more importantly, preserving charge neutrality in the device channel. The I_{on}/I_{off} well benchmarks with the record-high values already reported in literature

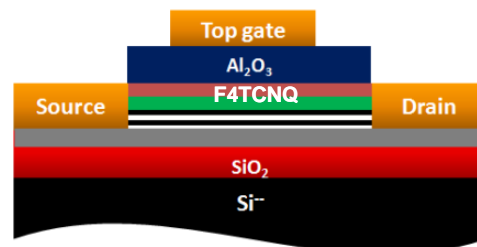


Fig. 1 Schematic of BLG-FET structure studied in this work.

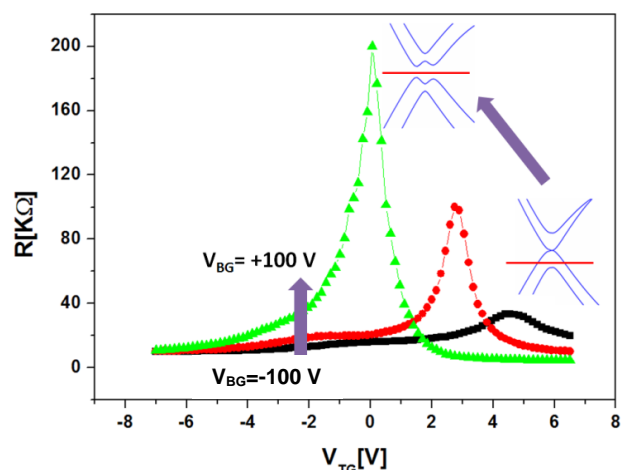


Fig. 2 Room temperature transfer characteristics (R vs. V_{TG}) of the BLG-FET described in this work. V_{BG} is varied from -100 to +100 V. The black curve corresponds to $V_{BG}=0$ V. The neutrality point voltage V_{NP} towards positive values is due to the effect of F₄TCNQ functionalization. The green and black curves are measured at $V_{BG}=+100 \text{ V}$ and $V_{BG}=-100 \text{ V}$, respectively.

References:

- [1] T. Ohta, A. Bostwick, T. Seyller, K. Horn and E. Rotenberg, *Science* **313**, 951 (2006)
- [2] F.N. Xia, D.B. Farmer, Y.M. Lin and P. Avouris, *Nano Lett* **10**, 715 (2010)
- [3] C.-T.L. Wenjing Zhang, Keng-Ku Liu, Teddy Tite, Ching-Yuan Su, Chung-Huai Chang, Yi-Hsien Lee, Chih-Wei Chu, Kung-Hwa Wei, Jer-Lai Kuo, and Lain-Jong Li, *Acs Nano* **5**, 7517 (2011)
- [4] C. Coletti, C. Riedl, D.S. Lee, B. Krauss, L. Patthey, K. von Klitzing, J.H. Smet and U. Starke, *Phys Rev B* **81** (2010)